

H_0 tension: shimmer of an unstable de Sitter cosmology?

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Ó Colgáin, van Putten, Yavartanoo, arXiv:1807.07451 (2018)

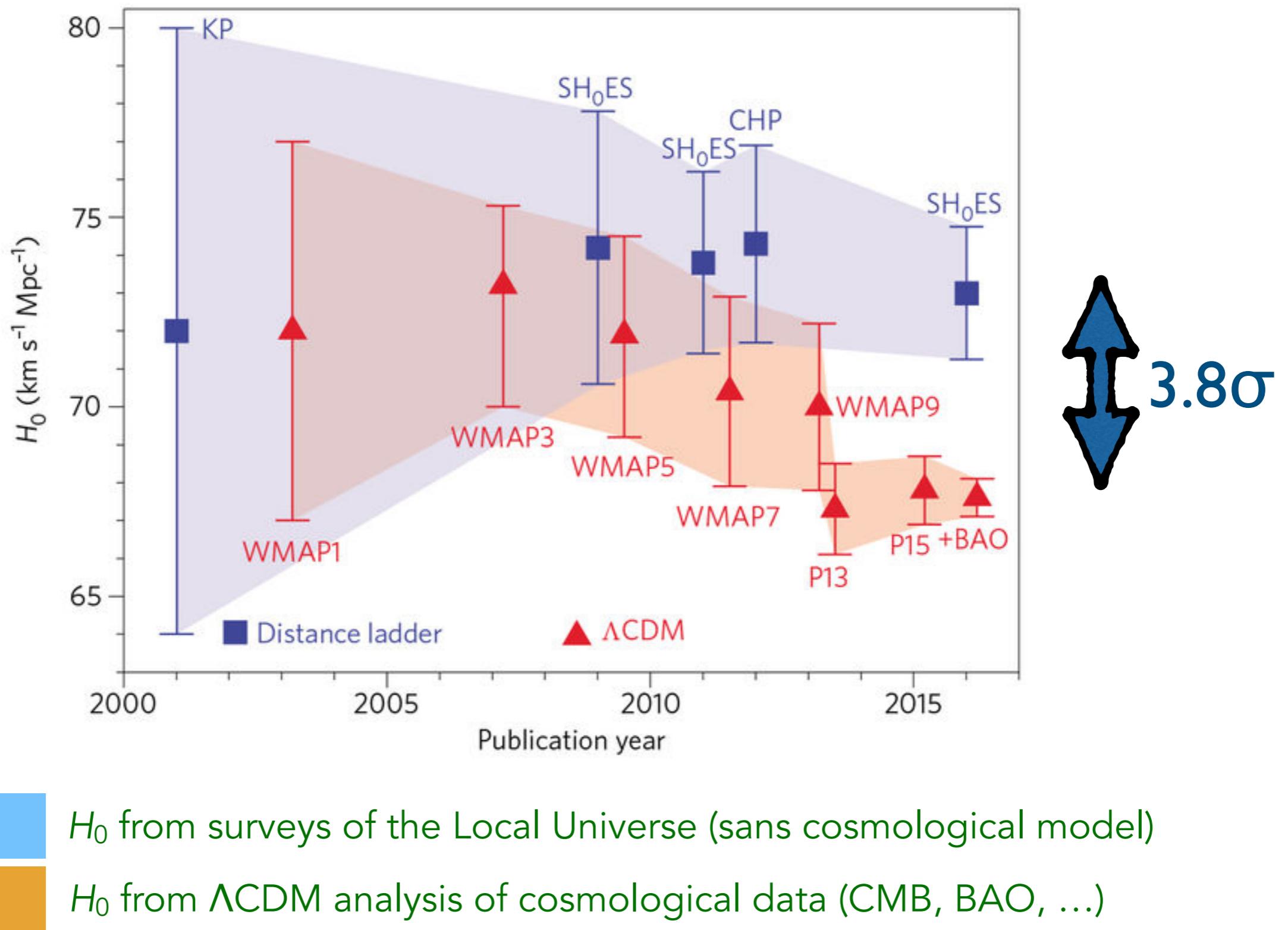
van Putten, 2017, ApJ **428**, 28

APCTP H_0 tension and Swampland: theory confronts reality

Dec 29 2018



H_0 tension

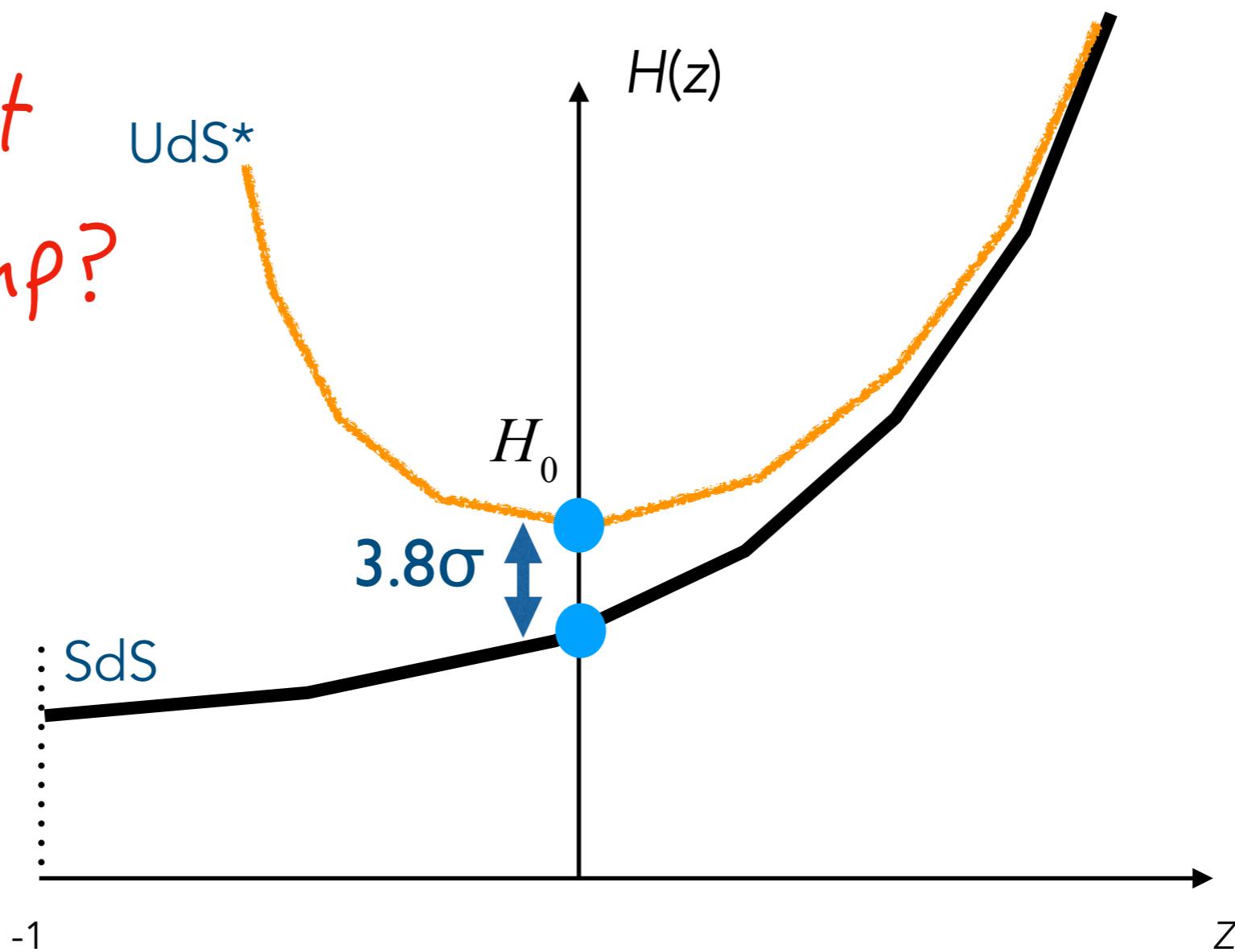


H_0 from surveys of the Local Universe (sans cosmological model)

H_0 from Λ CDM analysis of cosmological data (CMB, BAO, ...)

H_0 tension: shimmer from the future?

Future is out
of the Swamp?



UdS^* : Λ CDM ruled out in the future (*Swampland conjecture*),
 Λ CDM cannot hold to all orders today,
tensions today

*Vafa, C., 2005, hep-th/0509212

Brennan, T.D., Carta, F., and Vafa, C., 2017, arXiv:1711.00864;

Obied, Ooguri, Spodyneiko, Vafa, 2018

O'Colgain, E., van Putten, M.H.P.M., & Yavar, H., 2018, arXiv:1808.07451

Phase space - geometric scaling

Phase space BH spacetime is geometrically 2D (Bekenstein 1973 't Hooft 1993)

FRW-cosmologies have a cosmological horizon \mathcal{H} : an apparent horizon at Hubble radius $R_H=c/H$.

Turning this inside-out: phase space cosmological vacuum is 2D

This \mathcal{H} -constraint is woefully violated by Zel'dovich' UV-divergent dark energy **volume** density integral

Two polarizations:

$$M = 2\left(\frac{1}{2}T\right)S = \frac{H}{2\pi}\pi R_H^2 = \frac{1}{2}R_H$$

Gibbons & Hawking 1977
Verlinde 2011

Dark energy density:

Easson, Frampton & Smoot 2011

$$\rho_\Lambda = \frac{M}{V_H} = \frac{\frac{1}{2}R_H}{\frac{4\pi}{3}R_H^3} = \frac{3H^2}{8\pi} \equiv \rho_c$$

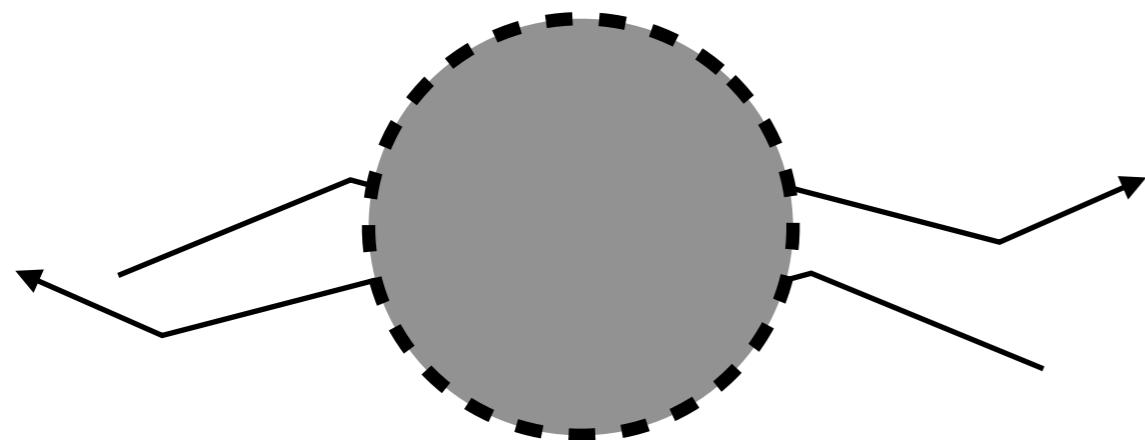
dS : residual of \mathcal{H} -constrained fluctuations

Is dS stable?

Apparent horizons: chromatic aberration

AH defined in the general relativity: geometric optics limit of gravitation

Subject to dispersion in IR:



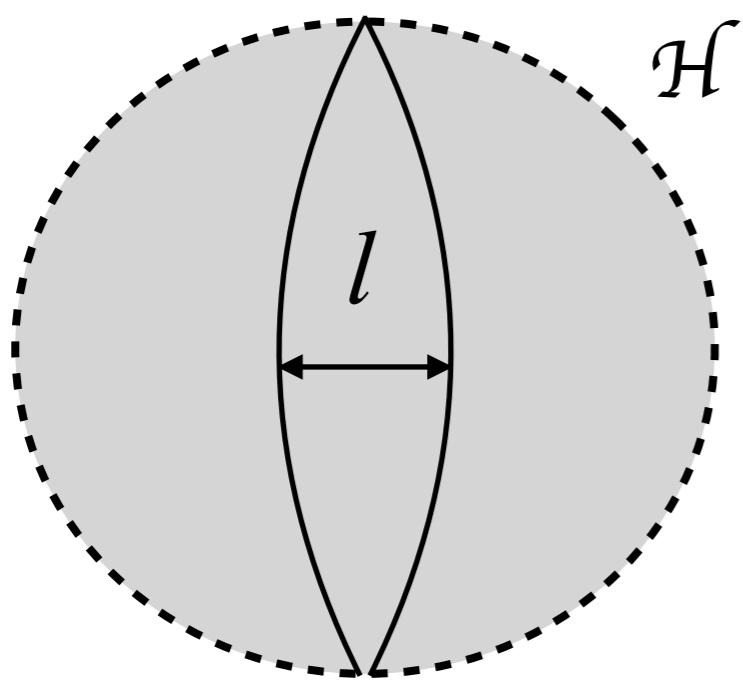
Hawking radiation, to null-infinity in
Minkowski spacetime

On-shell

Super-horizon scale fluctuations
leaking out

Super-horizon scale modes start leaking out at

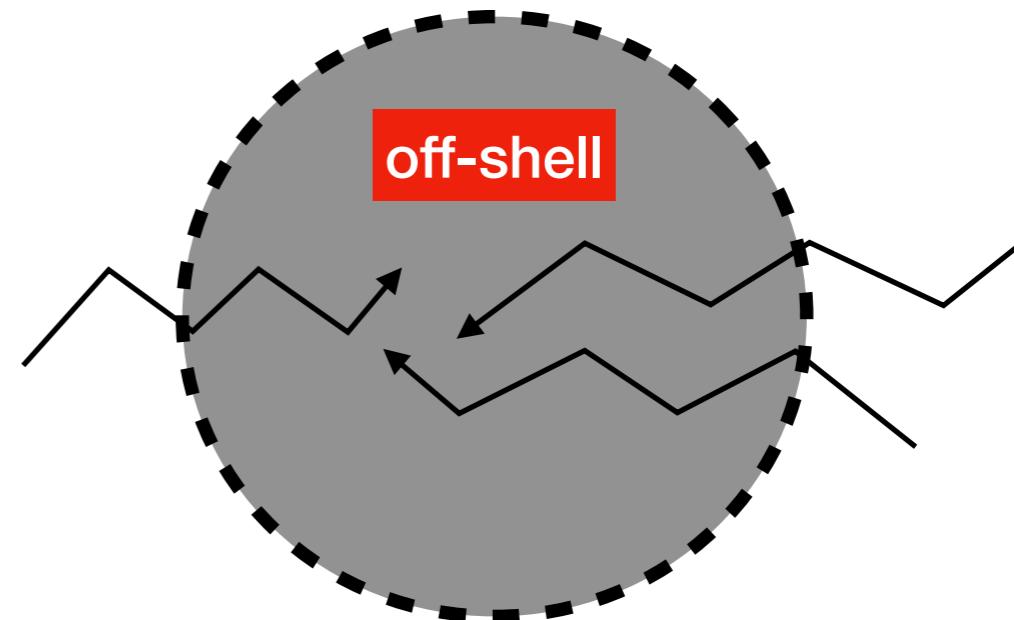
$$\omega_0 = \frac{1}{4\pi R_S}$$



*Following Gauss-Bonnet: normal
from tangential acceleration*

Applied to \mathcal{H} in FRW:

$$\omega_0 = \sqrt{1 - q} H$$



Radiation dominated ($q = 1$): $\omega_0 = 0$

Pick up $p < 0$ by imaginary wave numbers at super-horizon scale
fluctuations leaking in

Cosmological vacuum: dispersion

$$dS: \quad \omega_0 = \sqrt{2}H \quad (q = -1)$$

Wave equations of EM/U(1) and GW/SO(3,1) in Lorenz gauge (van Putten & Eardley 1996): coupling to Ricci tension picks up sensitivity to Λ

Dispersion relation of linearized wave-motion: $\omega = \sqrt{k^2 + \Lambda}$

$$\Lambda = \omega_0^2 \quad \text{from } \mathcal{H}$$

Quadratic: inherently small and positive today

van Putten 2017 ApJ 837 22

Equations of motion

Three-flat FRW cosmology:

$$G_{ab} = 8\pi T_{ab} + (1 - q)H^2 g_{ab}$$

Density parameters

$$\Omega_\Lambda = \frac{1}{3}(1 - q)$$

$$\Omega_M = \frac{1}{3}(2 + q)$$

Dark “matter”

$$q = 3\Omega_p$$

$$w = \frac{2q - 1}{1 - q}$$

Radiation picture

$$T_{ab} = (\rho + p)u_a u_b + pg_{ab} - \omega_0^2 g_{ab}$$

$$T_{ab} = \rho_c \left[(1 - q)\pi_{ab}^- + q\pi_{ab}^+ \right]$$

$$\pi_{ab}^\pm = \text{dia}(1, \pm \frac{1}{3}, \pm \frac{1}{3}, \pm \frac{1}{3})$$

- * Radiation dominated ($q=1$)
- * Matter dominated ($w=0, q=1/2$)
- * de Sitter ($q=-1$)

Analytic solution

$$H(z) = H_0 \sqrt{1 + \omega_m \left(6z + 12z^2 + 12z^3 + 6z^4 + \frac{6}{5}z^5 \right)} (1+z)^{-1}$$

van Putten, 2017, ApJ, 848, 28

Analytic solution

$$H(z) = H_0 \sqrt{1 + \omega_m \left(6z + 12z^2 + 12z^3 + 6z^4 + \frac{6}{5}z^5 \right)} (1+z)^{-1}$$

van Putten, 2017, ApJ, 848, 28

Same parameters of late-time Λ CDM cosmology:

$$H_0 = H(0), \quad \omega_m = \Omega_m(0)$$

Analytic solution

$$H(z) = H_0 \sqrt{1 + \omega_m \left(6z + 12z^2 + 12z^3 + 6z^4 + \frac{6}{5}z^5 \right)} (1+z)^{-1}$$

van Putten, 2017, ApJ, 848, 28

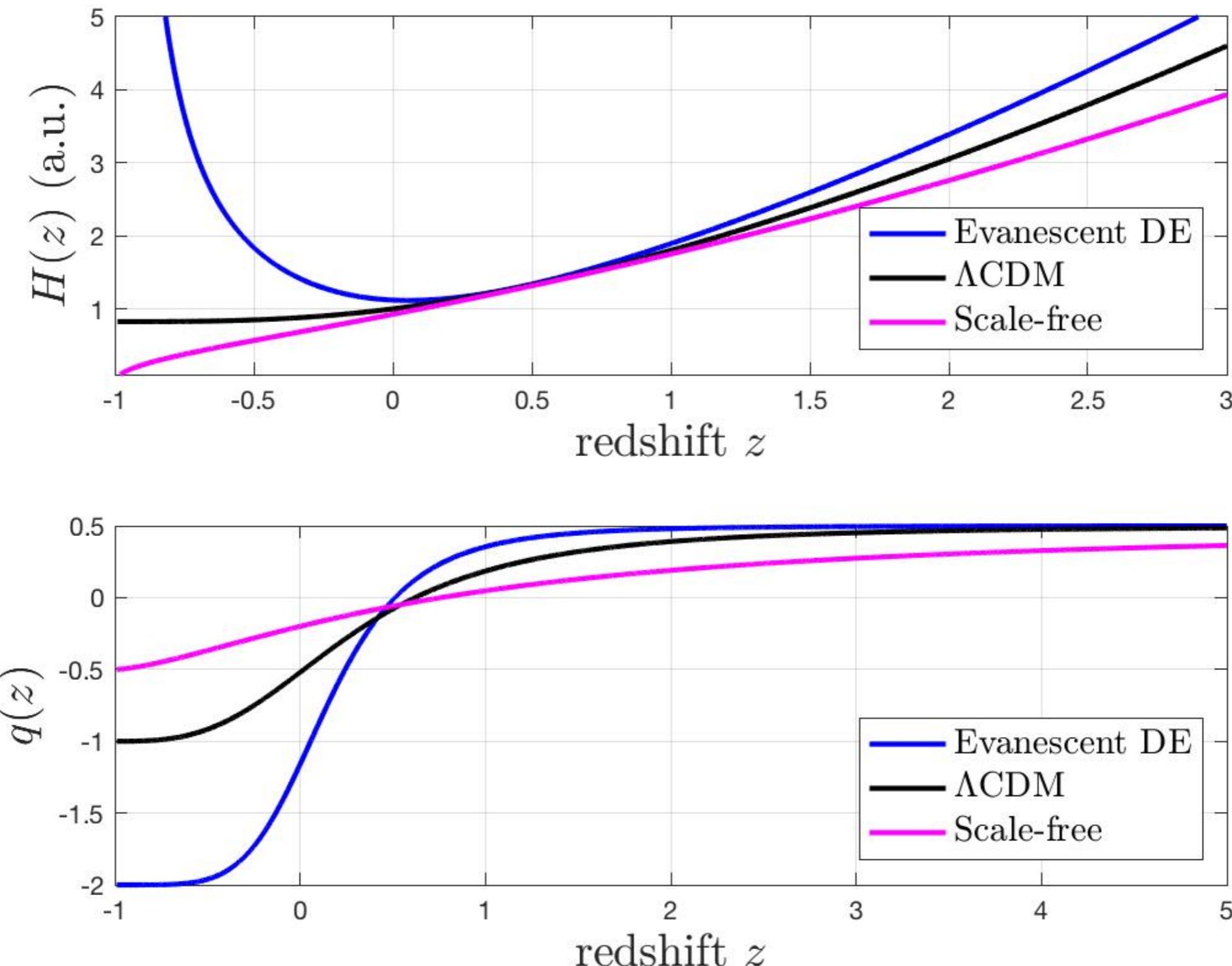
Same parameters of late-time Λ CDM cosmology:

$$H_0 = H(0), \quad \omega_m = \Omega_m(0)$$

Out-of-Swampland:

UdS: $H(z) \rightarrow \infty \quad (z \rightarrow -1)$

Analytic solution



H_0 tension: observational test

“Swampland escape” at nearby *turning point* in Hubble parameter

$$z_* = \left(\frac{5 - 6\omega_m}{9\omega_m} \right)^{\frac{1}{5}} - 1 \simeq 0.07 \quad (\omega_m = 0.28)$$



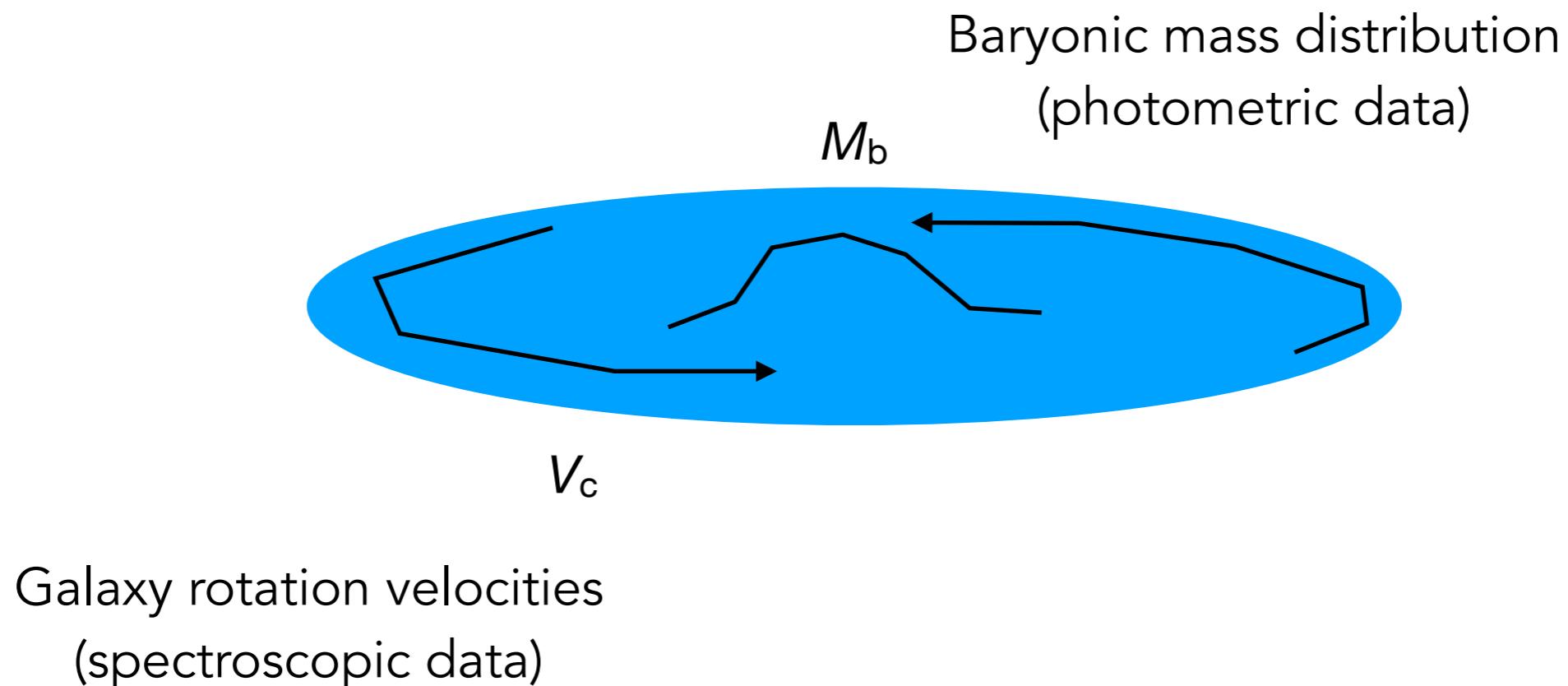
1 dimensionless parameter amenable to measurement by a single (homogeneous) survey of the Local Universe

Ó Colgáin, van Putten, Yavartanoo, arXiv:1807.07451 (2018)

What about dark matter?

Three-flatness requires dark/phantom “stuff” on cosmological scales

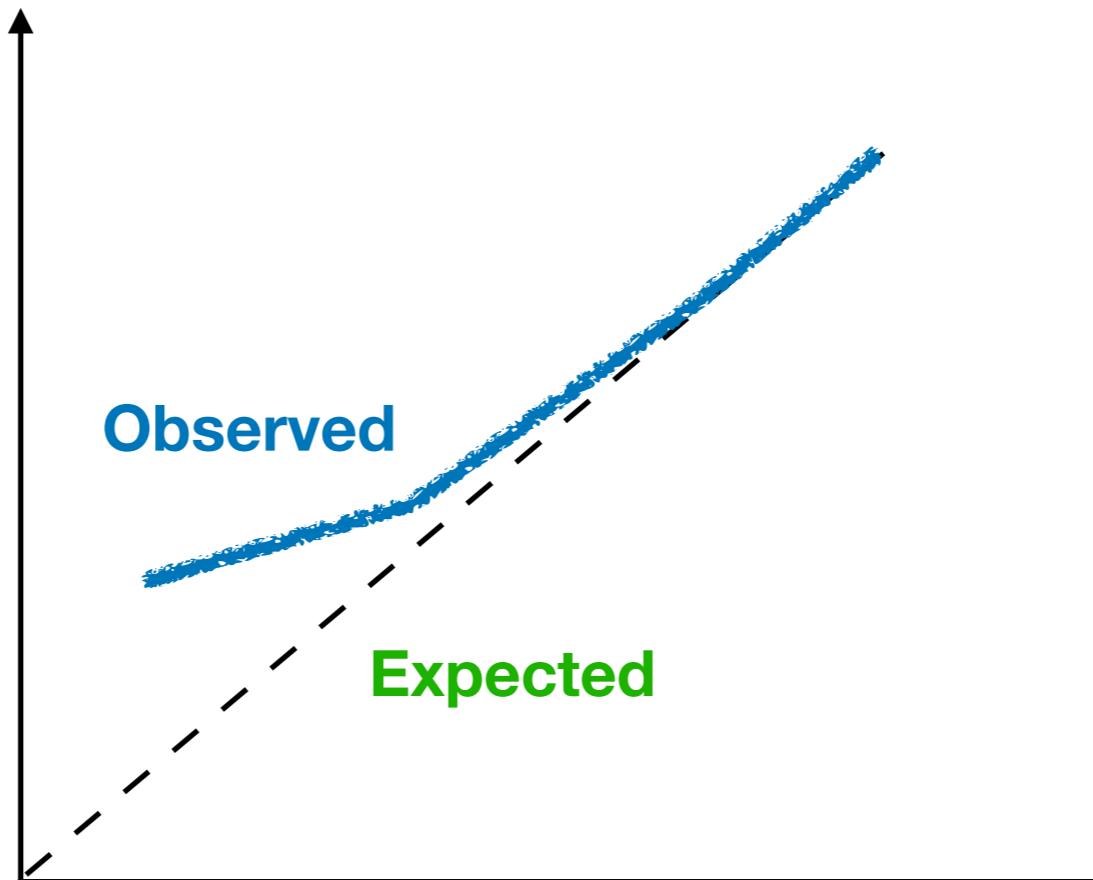
On galactic scales, data tell a different story...



Galaxy rotation curves

Observed/Expected
radial acceleration

$$\frac{\alpha}{a_N} = \frac{V_c^2 / r}{V_b^2 / r}$$

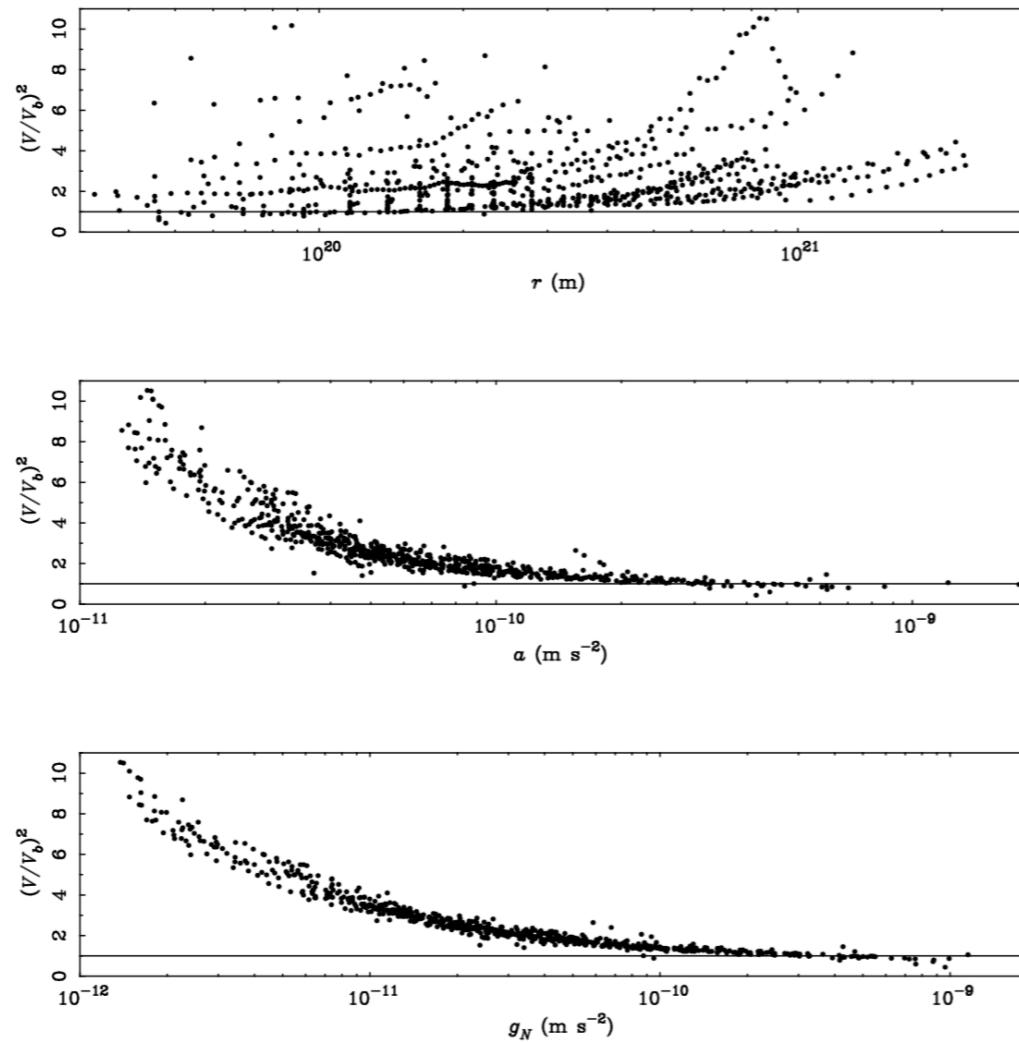


$$a_N \equiv -\frac{GM_b}{r^2}$$

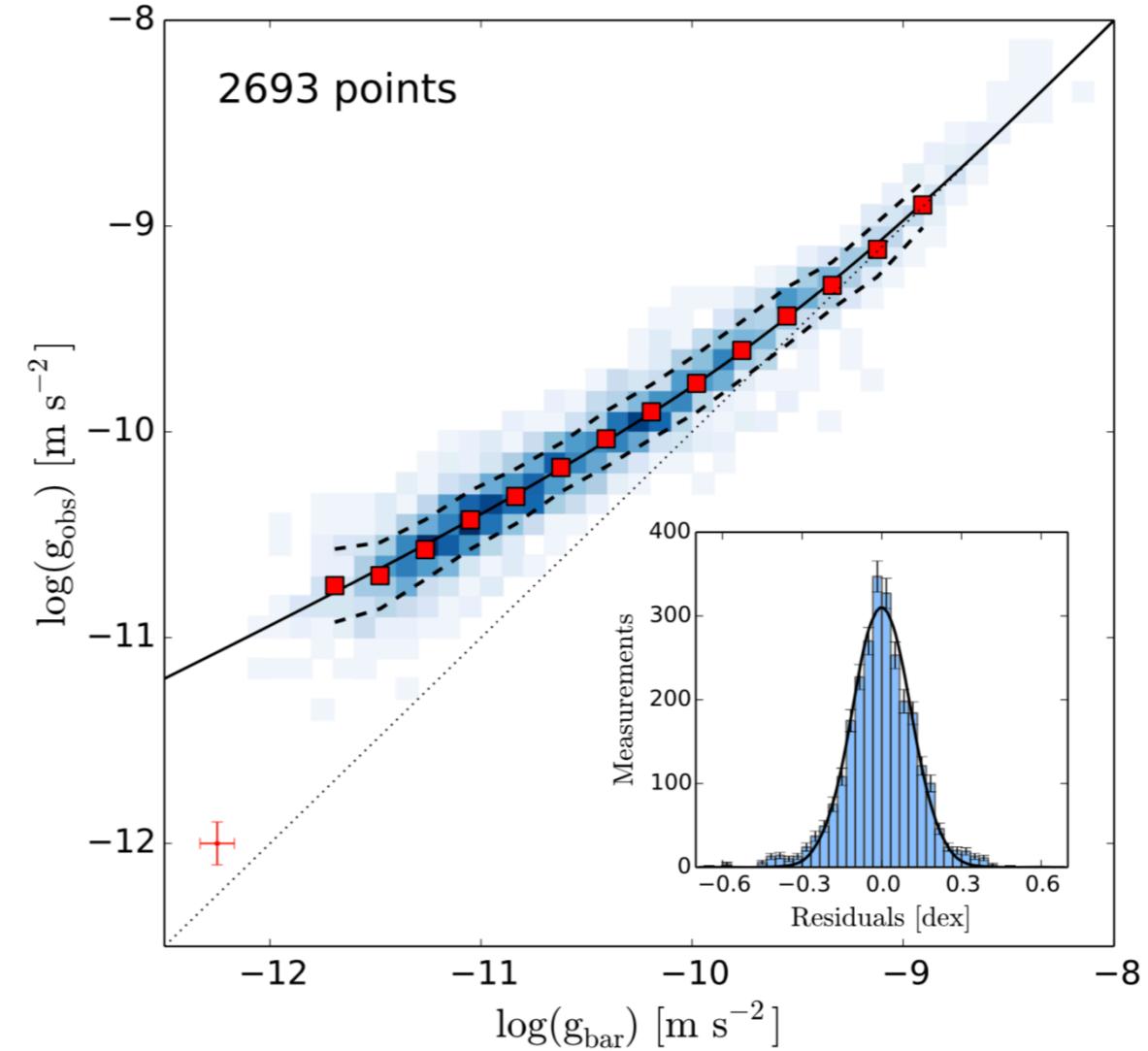
Famae, B., & McGaugh, S.S., 2012, LRR, 15, 10

McGaugh, S.S., Lelli, F., & Shomber, J., 2016, PRL, 117, 201101

Galaxy rotation curves



SPARC: Spitzer Photometry and Accurate Rotation curves Catalogue (z=0)



Famae, B., & McGaugh, S.S., 2012, LRR, 15, 10

McGaugh, S.S., Lelli, F., & Shombert, J., 2016, PRL, 117, 201101

Galaxy rotation curves - normalized to a_{dS}

$$a_{dS} \equiv cH$$

Photometry

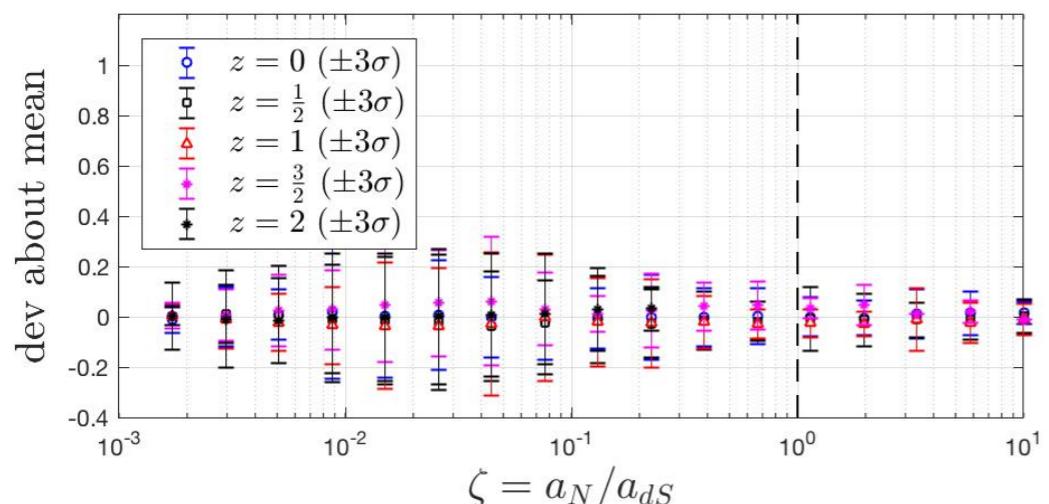
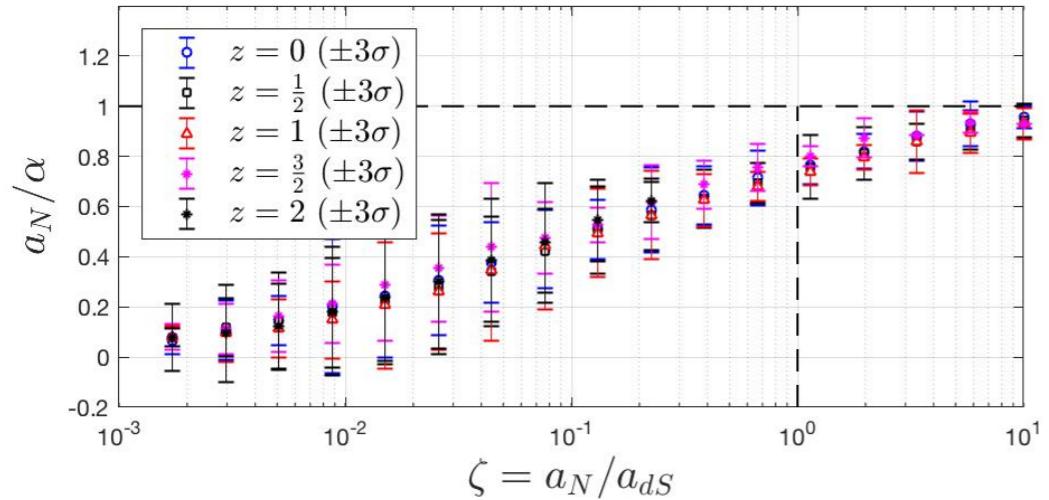
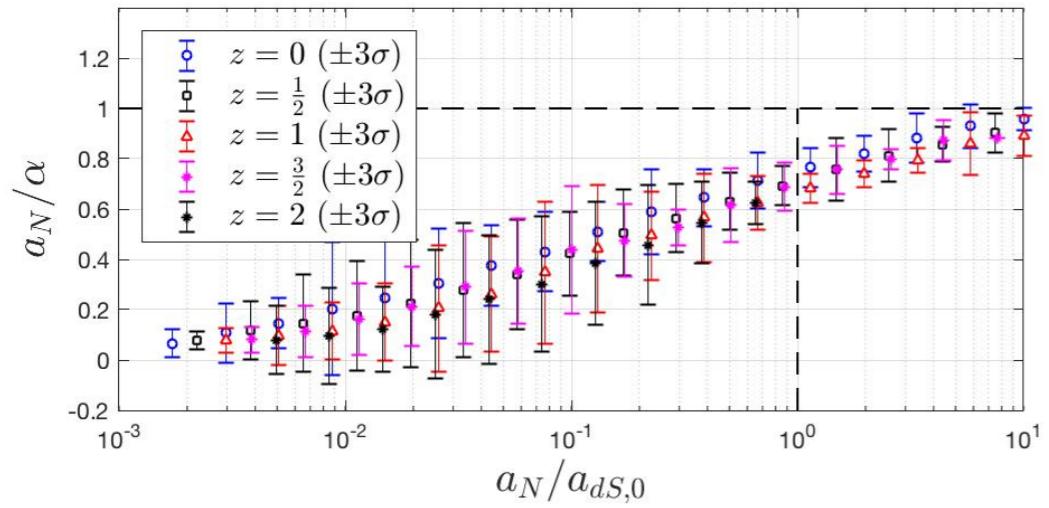
$$\zeta = \frac{a_N}{a_{dS}}, \quad a_{dS} \equiv cH$$

Spectroscopy

Radial acceleration⁻¹

$$\frac{a_N}{\alpha}$$

van Putten, 2018, MNRAS 481 L26



McMaster Unbiased Galaxy Simulations 2

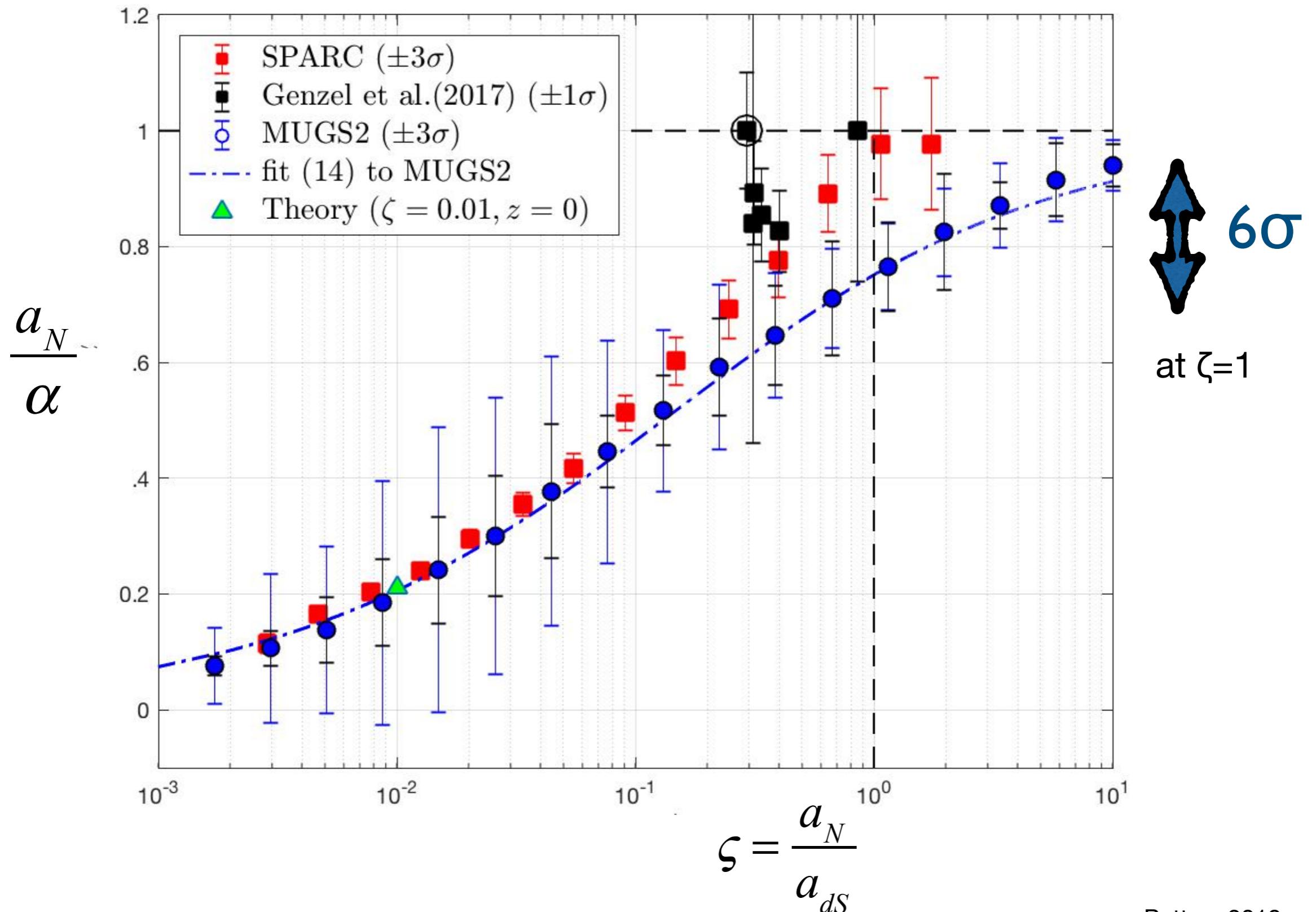
Keller, B.W., Wadsley, J., Benincasa, S.M., & Couchmanm,
H.M.P., 2014, MNRAS, 442, 3013

Keller, B.W., Wadsley, J., & Couchman, H.M.P., 2016,
MNRAS, 463, 1431

Keller, B.W., & Wadsley, J.W., 2017, ApJ, 835, L17

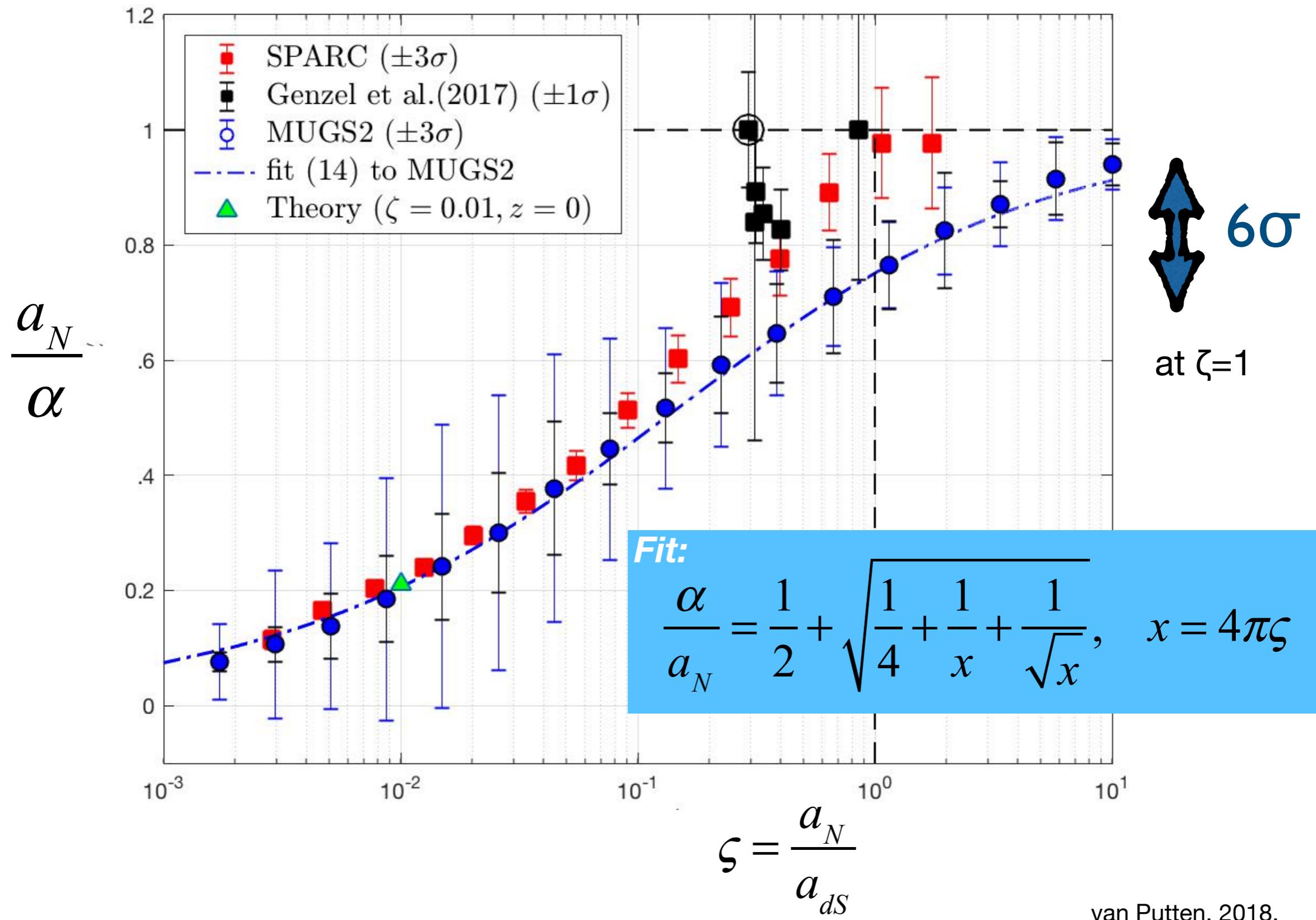
Self-similar galaxy dynamics
tracing background cosmology

van Putten, 2018, MNRAS 481 L26

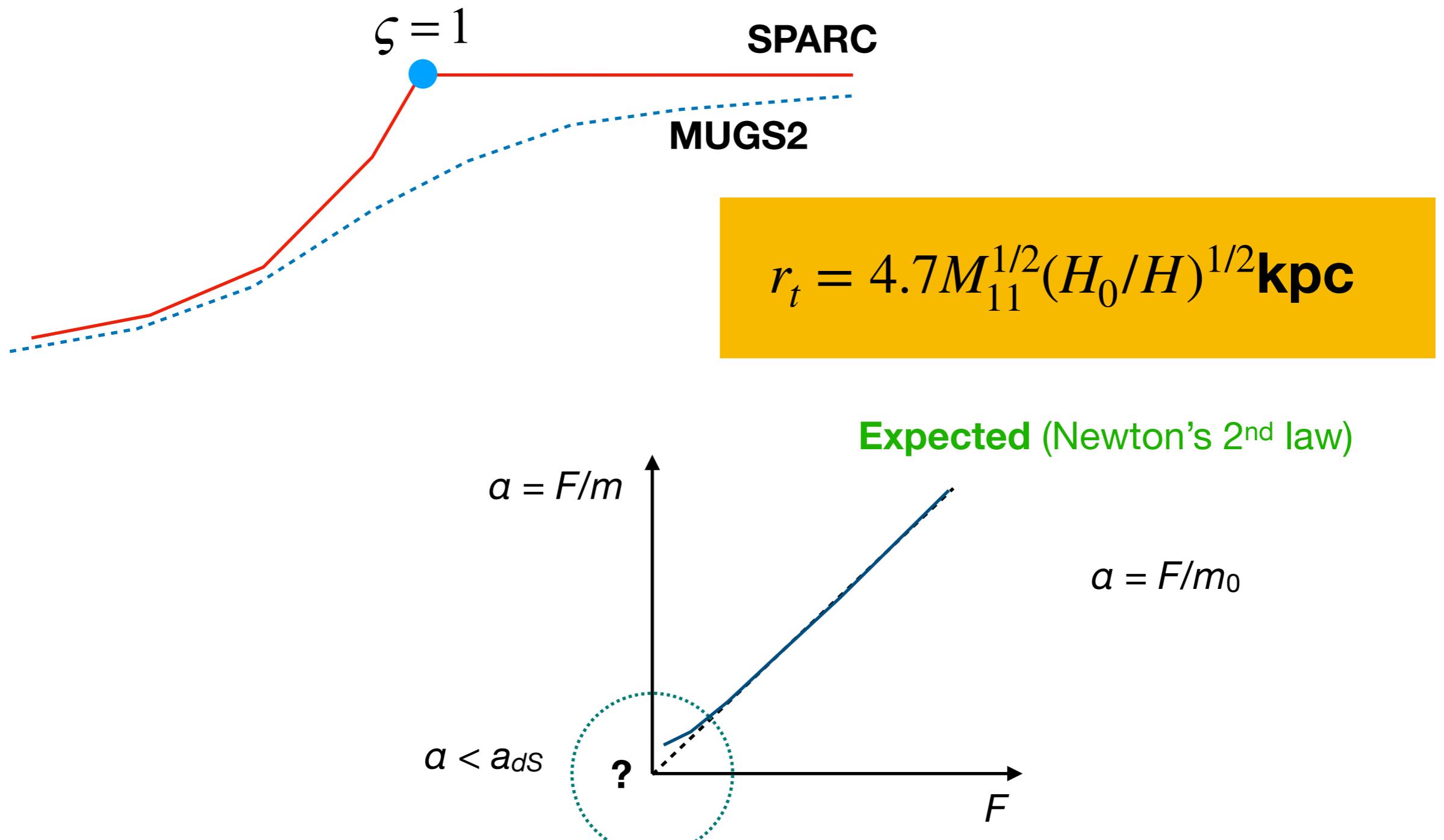


van Putten, 2018,
MNRAS, 481, L26

Λ CDM/MUGS2 vs MUGS2

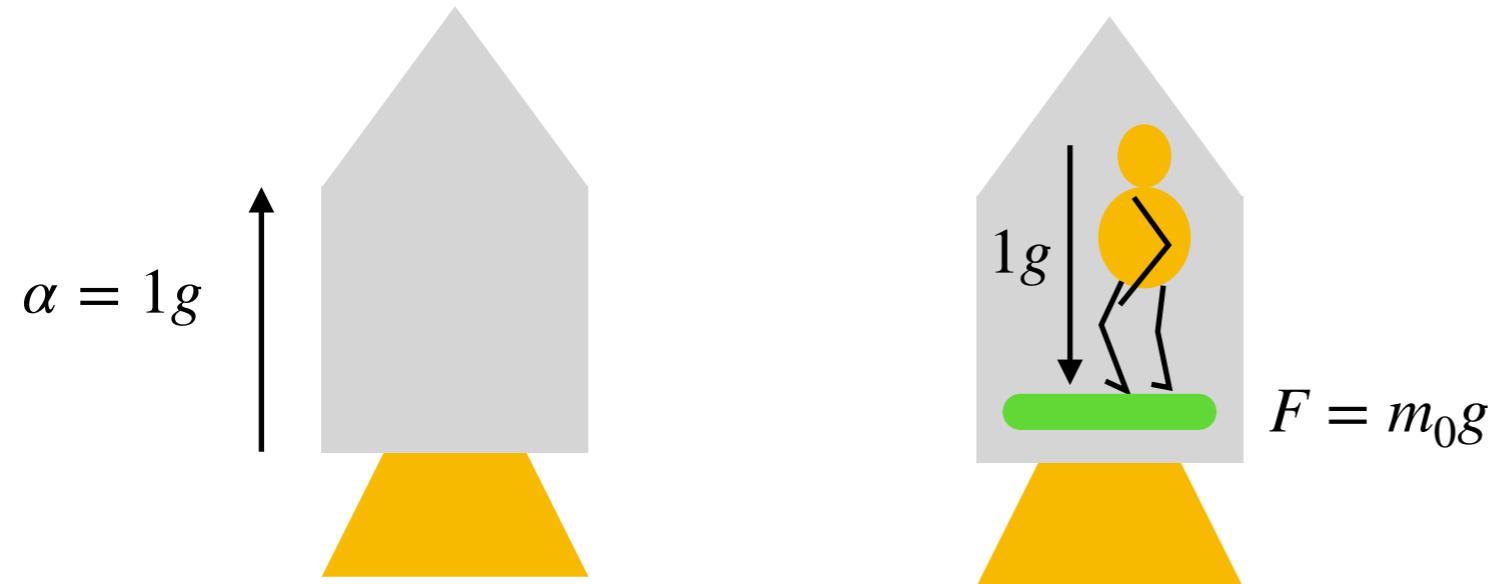
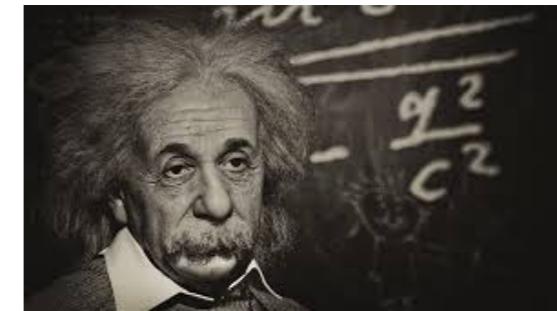


Approximate C^0 galaxy dynamics



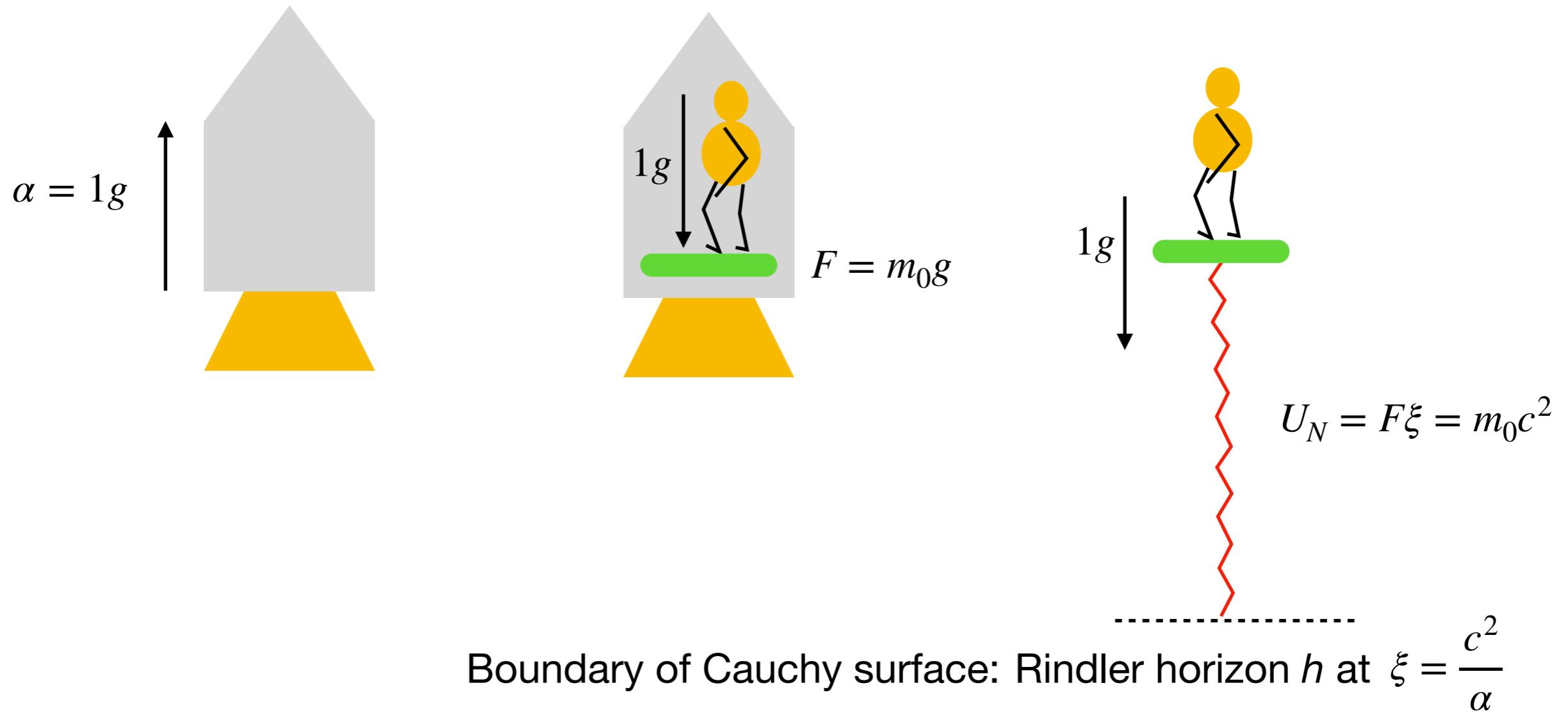
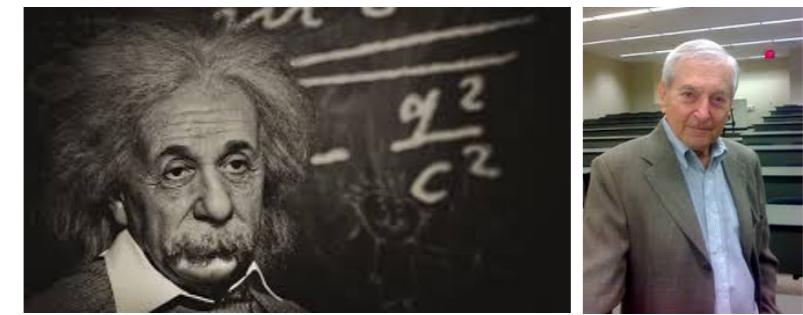
6σ gap: problem of inertia... ?

Equivalence Principle



Apples fall, light is bend, ... : gravitational field observed by the traveler inside

Extended Equivalence Principle



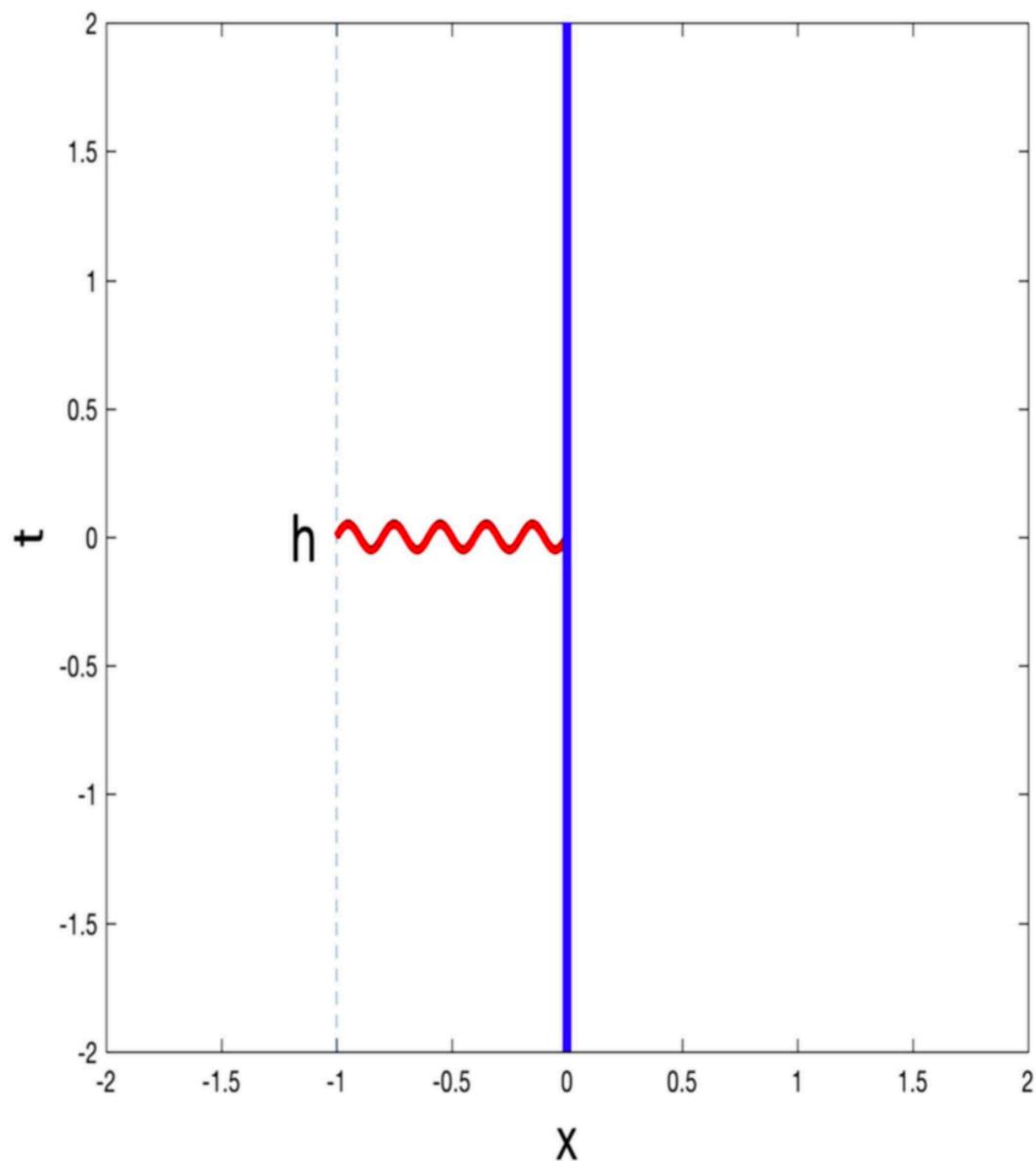
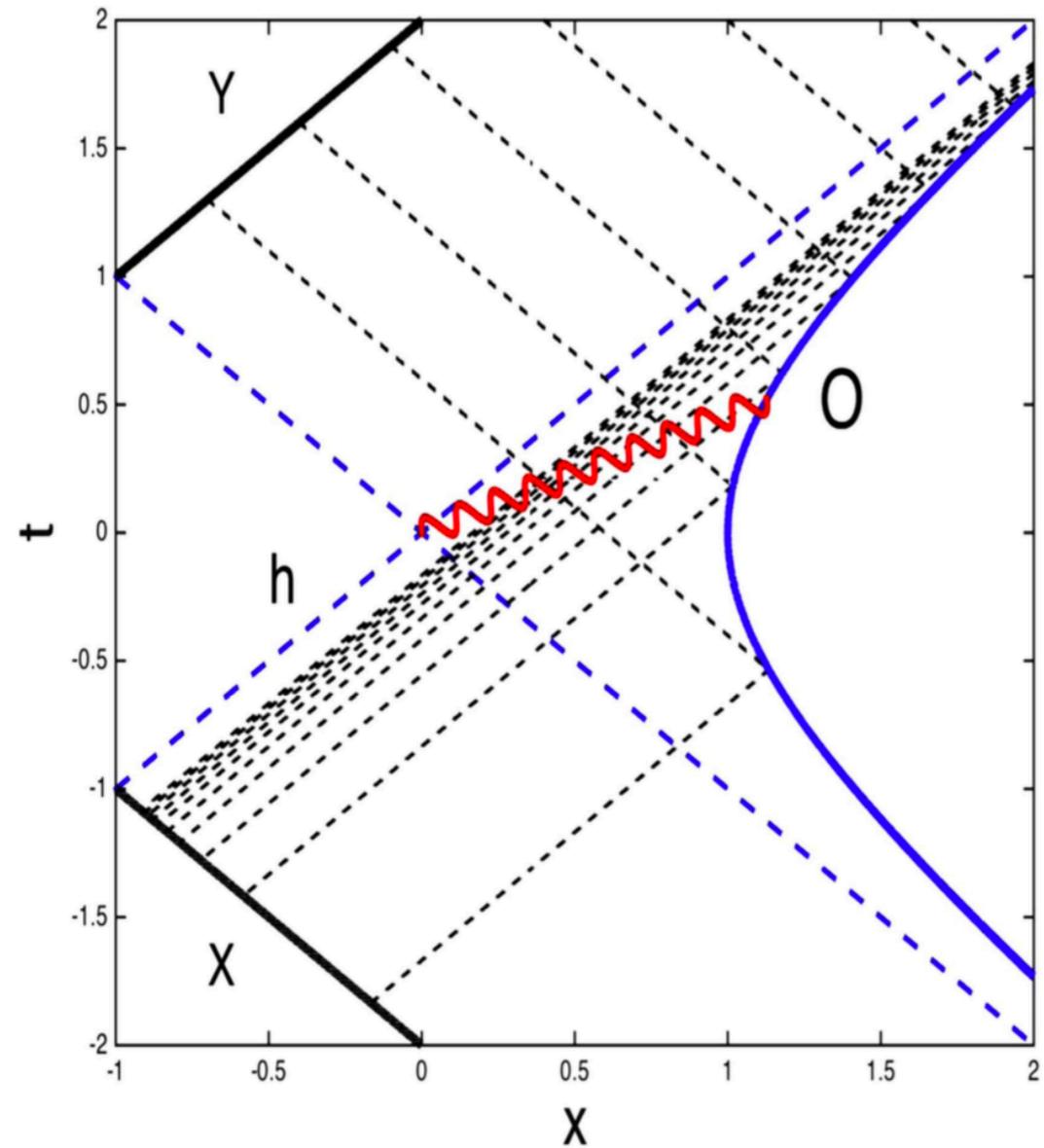
Newtonian inertia = Gravitational binding energy to h

van Putten, 2017, ApJ, 837, 22; ApJ, 848, 28

Inertia from entanglement entropy with h

THE ASTROPHYSICAL JOURNAL, 837:22 (8pp), 2017 March 1

van Putten

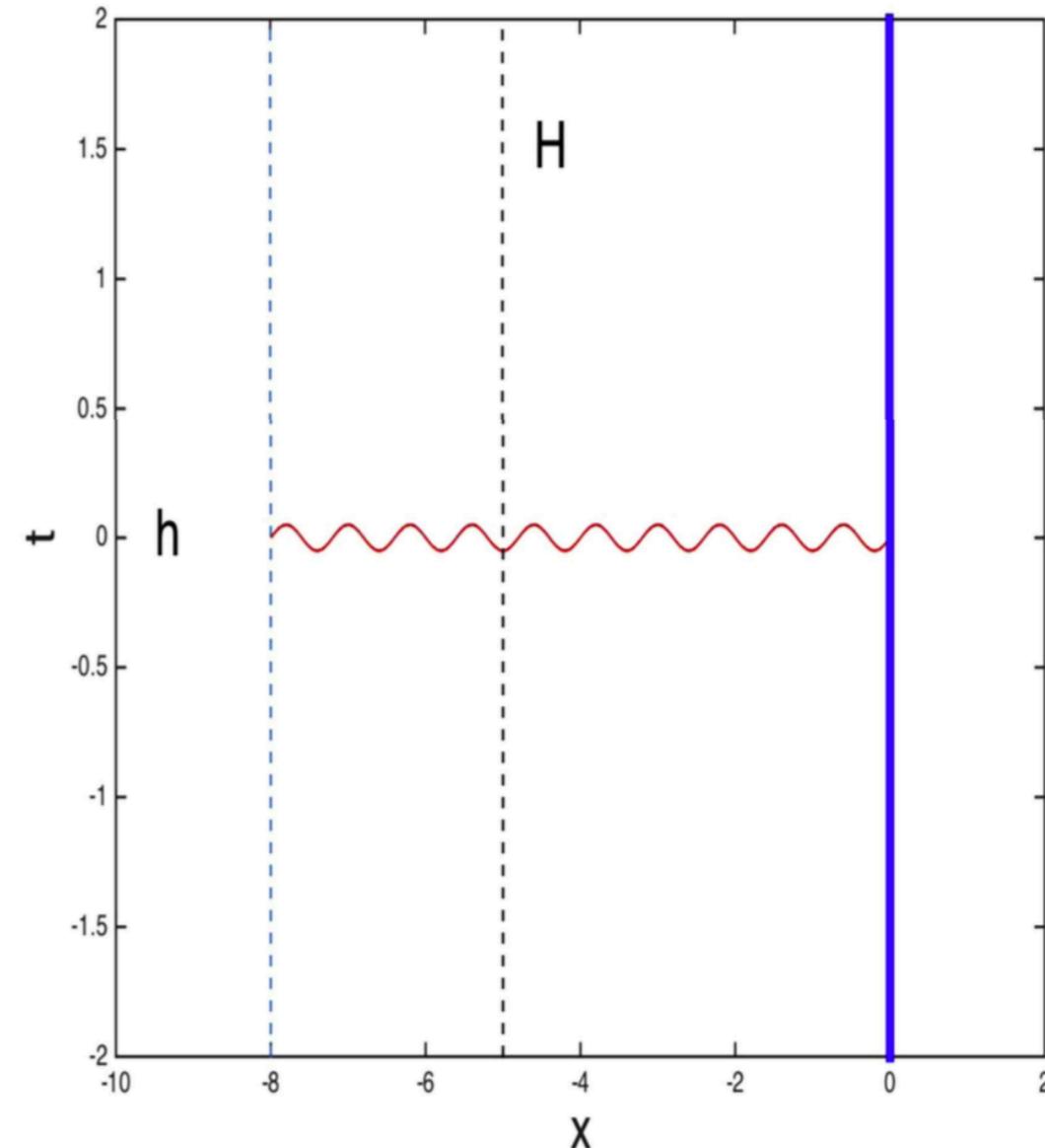
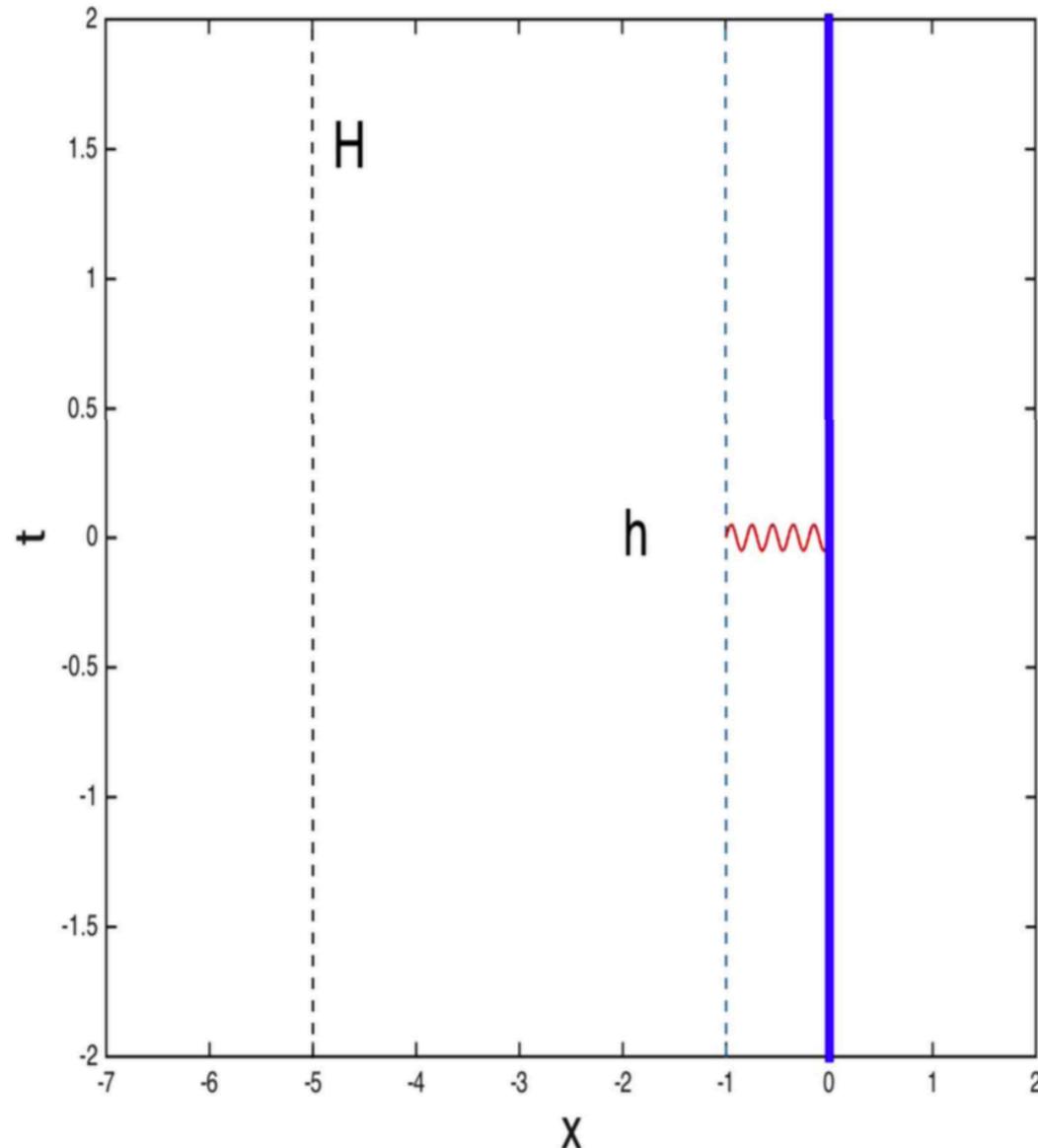


$$U_N = \int_0^\xi T_U dS = m_0 c^2, S = 2\pi m_0 \xi$$

van Putten, 2017, ApJ, 837, 22; ApJ, 848, 28

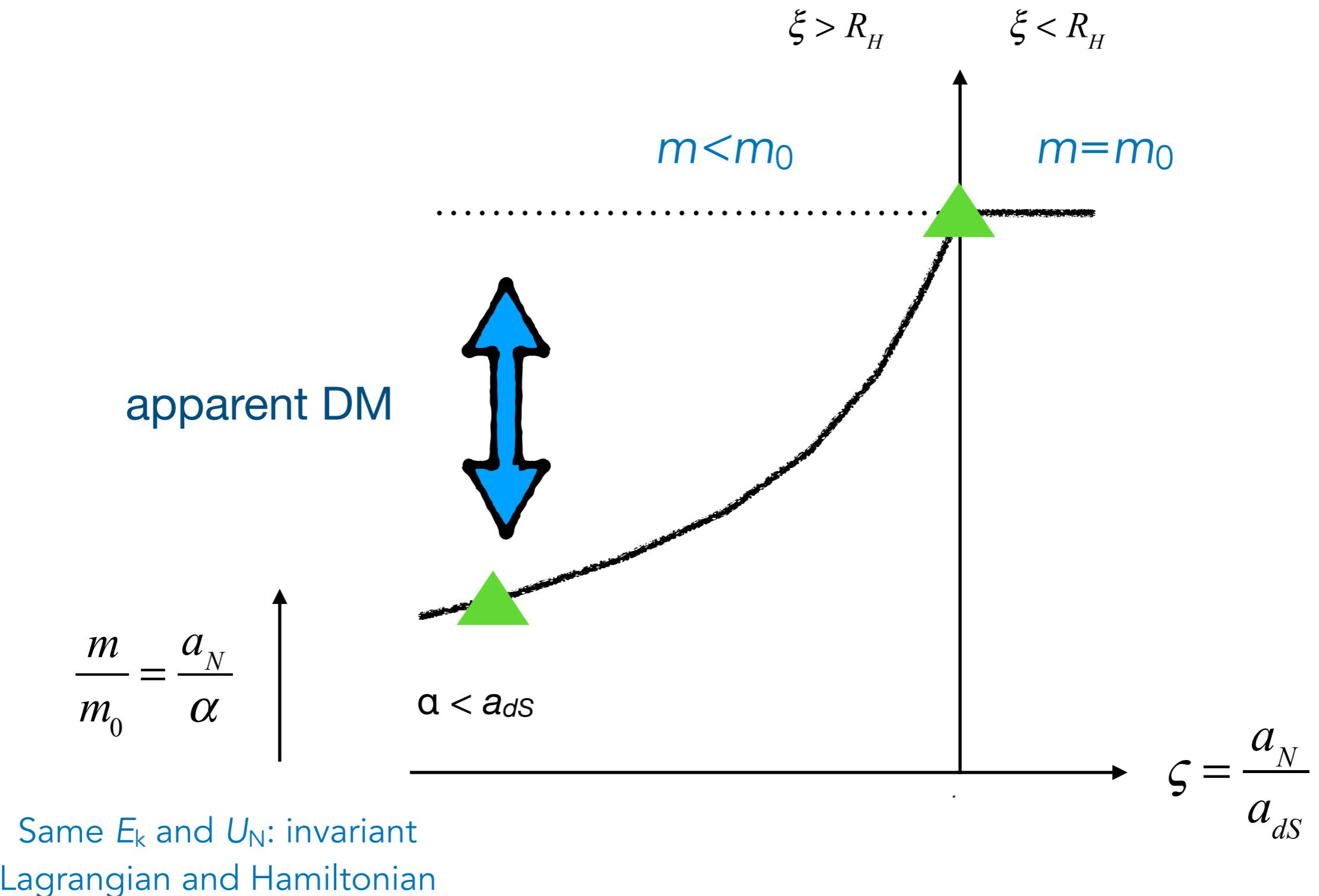
C^0 transition: entanglement to \mathcal{H}

$$R_H = \frac{c^2}{a_{dS}} = \frac{c}{H}$$



van Putten, 2017, ApJ, 837, 22; ApJ, 848, 28

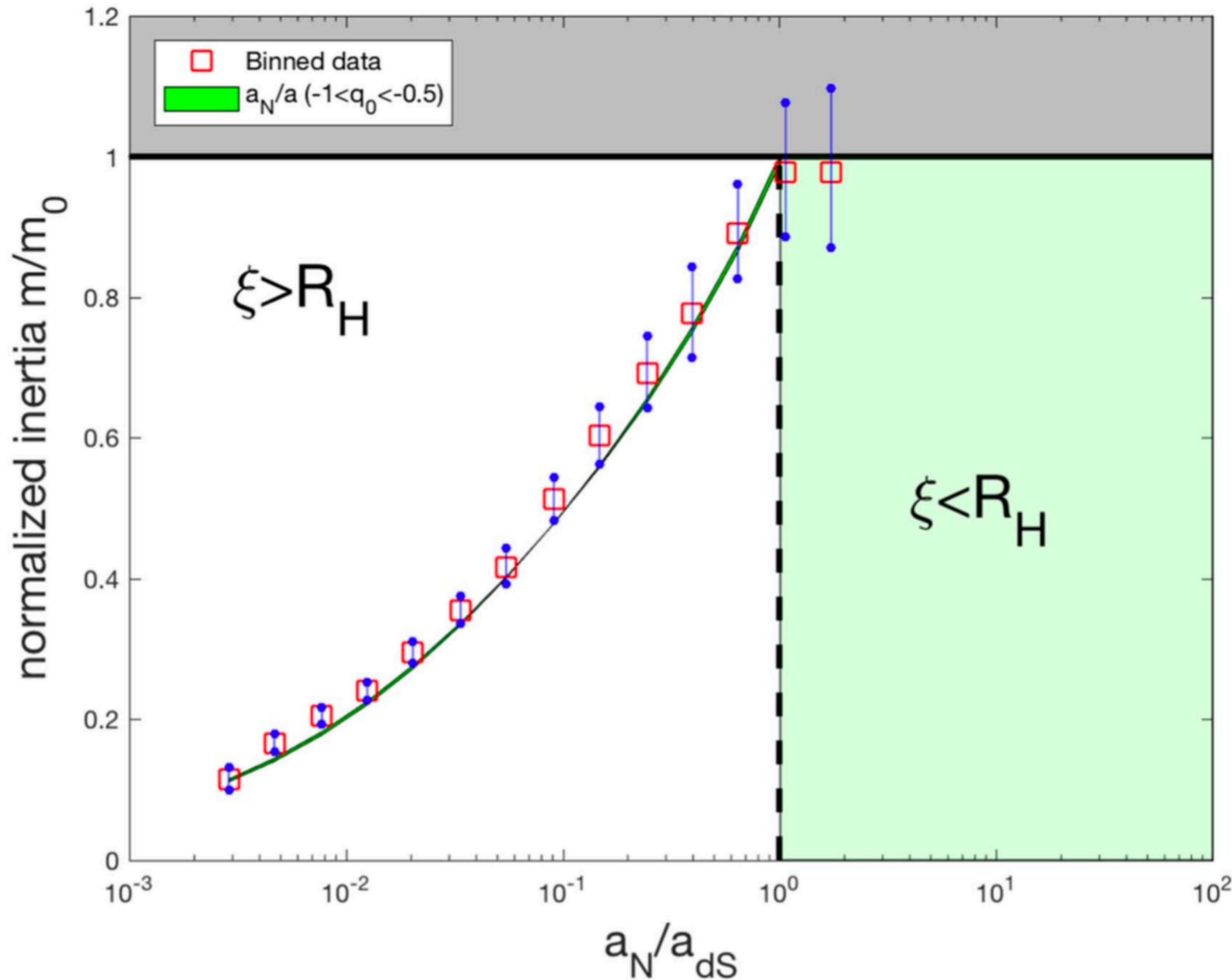
C^0 galaxy dynamics across a_{dS}



van Putten, 2017, ApJ, 848, 28

C^0 galaxy dynamics in SPARC

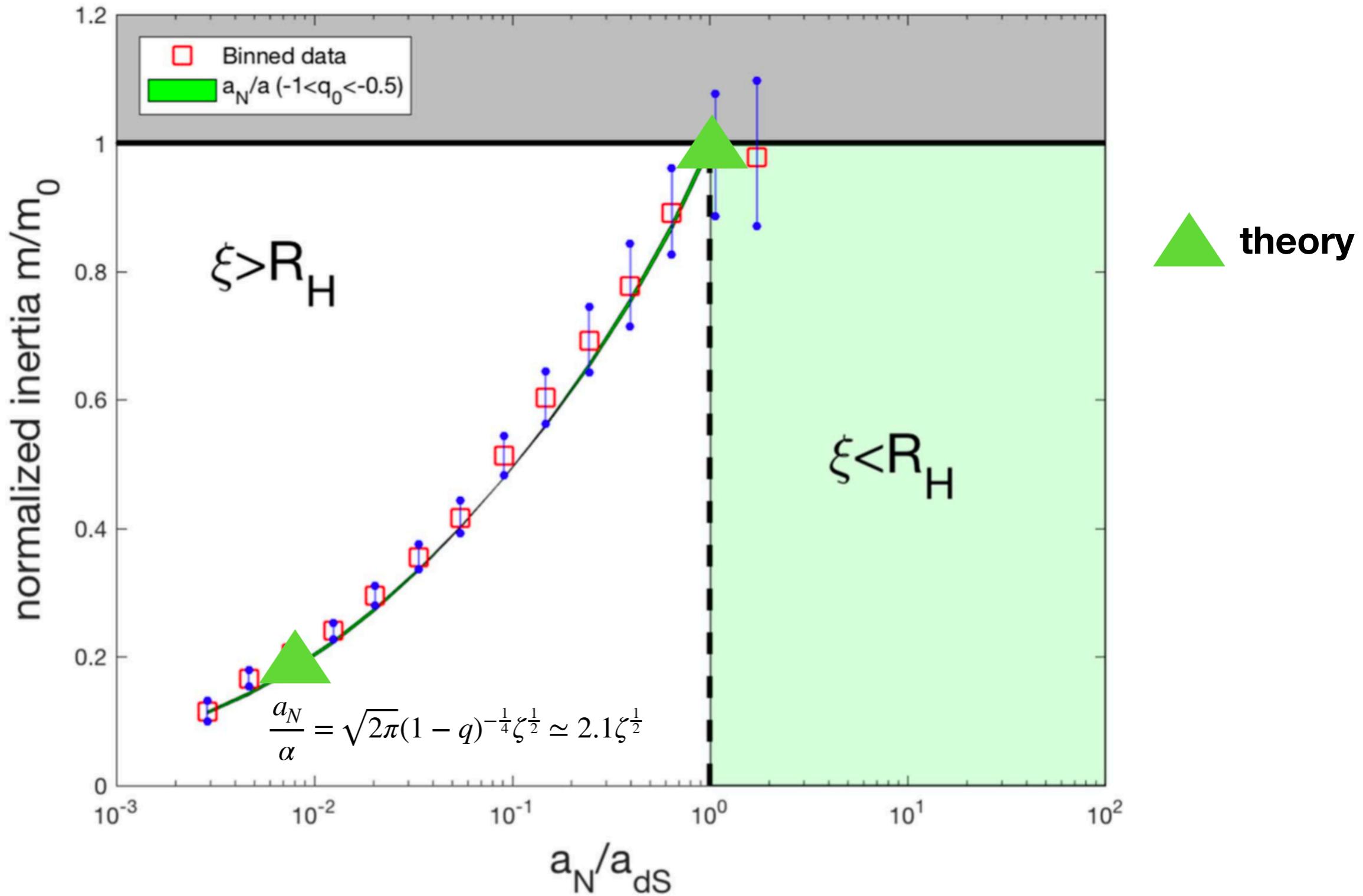
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van Putten, 2017, ApJ, 837, 22

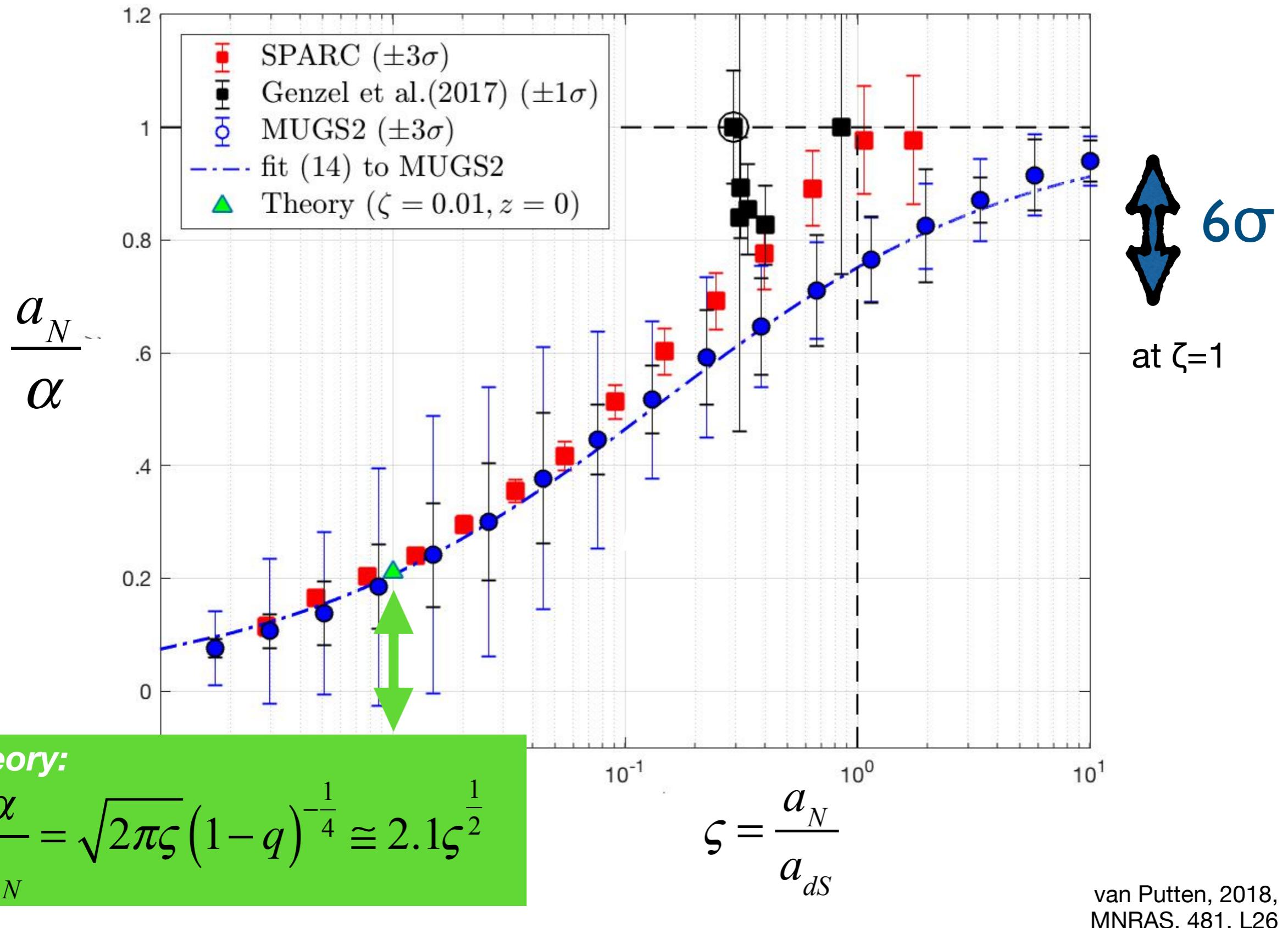
C^0 galaxy dynamics in SPARC

THE ASTROPHYSICAL JOURNAL, 837:22 (8pp), 2017 March 1



van Putten, 2017, ApJ, 837, 22

C^0 galaxy dynamics in SPARC



van Putten, 2018,
MNRAS, 481, L26

Conclusions and observational outlook

In Bekenstein-'t Hooft 2D scaling of phase space, \mathcal{H} protects vacuum against Zel'dovich UV catastrophe in Λ . **The residual is de Sitter.**

- \mathcal{H} in FRW implies $\Lambda = (1-q)H^2$ with dS unstable. Future is *out of Swampland*, ΛCDM cannot hold true to all orders today, anticipating tension in H_0 - **3.8 σ today**.
- For $H_0 \sim 73.5$ km/s/Mpc, it gives a test of “Swampland escape” by a **turning point** at

$$z_* = \left(\frac{5 - 6\omega_m}{9\omega_m} \right)^{\frac{1}{5}} - 1 \simeq 0.07 \quad (\omega_m = 0.28)$$

Dark “matter” is required for three-flatness but NOT for galaxies

- $\sim C^0$ galaxy dynamics/SPARC is not accounted for by ΛCDM : **6 σ failure.**
- May be due to Newton's 2nd law: inertia switching entanglement to \mathcal{H} at $\alpha < a_{\text{dS}}$:

$\lambda_C \gg r_t$: $m_{\text{DM}} \ll 10^{-27}\text{eV}$ outside realm of lab experiments.